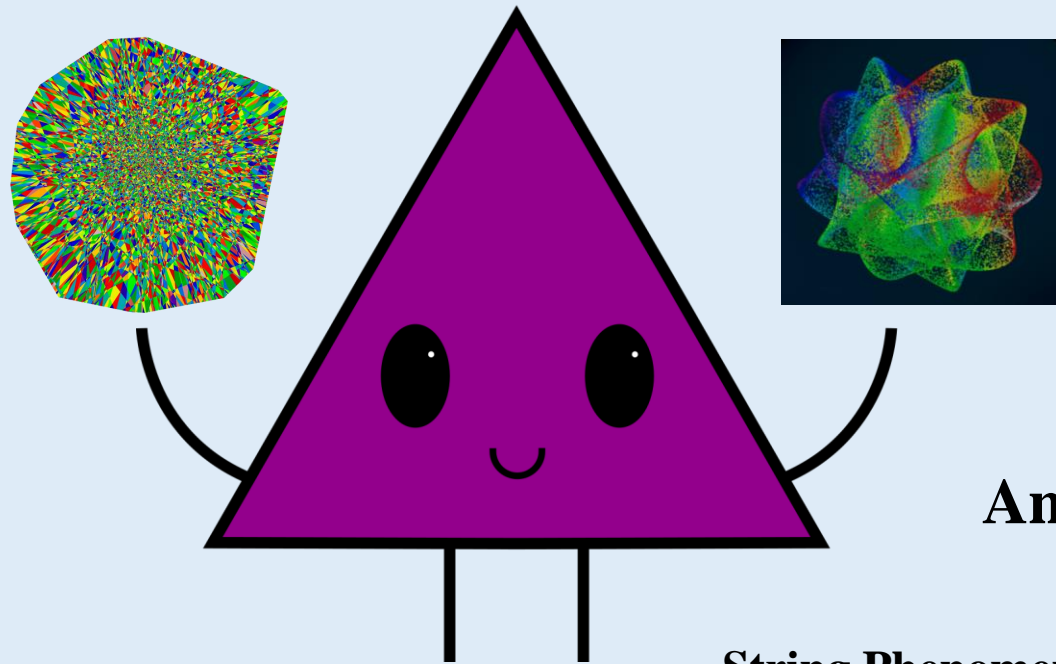


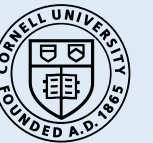
Convergence of Worldsheet Instanton Corrections in AdS Flux Vacua



Andres Rios-Tascon

Cornell University

String Phenomenology conference 2022



Based on work with M. Demirtas, M. Kim, L. McAllister, J. Moritz.

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Takeaways:

- Our improvements in computing Gopakumar-Vafa invariants allow for more complex constructions and robust checks of control (among many other things!).
- These tools will be integrated into our CYTools package, which will be released soon.

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Gukov-Vafa-Witten
flux superpotential



Non-perturbative contributions
from ED3-branes or strong gauge
dynamics on stacks of seven-
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We start by engineering $W_0 := \langle |W_{\text{flux}}| \rangle \ll 1$. We do this by picking fluxes that make the perturbative part vanish, so that only contribution are from IIA worldsheet instantons and leading terms form a racetrack. [Demirtas, Kim, McAllister, Moritz '19]

$$W_{\text{flux}}(\tau) = c(e^{2\pi i p_1 \tau} + A e^{2\pi i p_2 \tau}) + \dots$$

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In our flagship example we have

$$W_{\text{flux}}(\tau) \propto -2e^{2\pi i \tau \cdot \frac{7}{29}} + 252e^{2\pi i \tau \cdot \frac{7}{28}} + \mathcal{O}\left(e^{2\pi i \tau \cdot \frac{43}{116}}\right)$$

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Numbers in red are Gopakumar-Vafa (GV) invariants.

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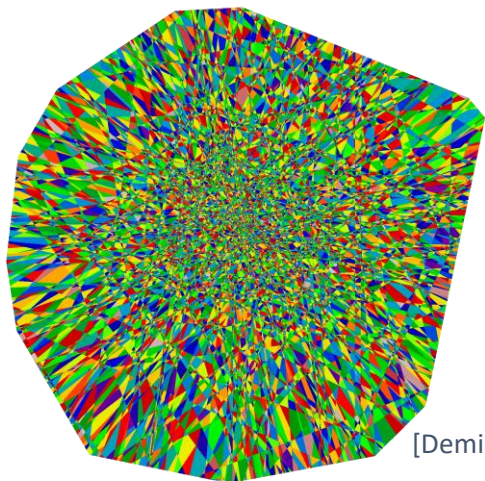
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[Demirtas, McAllister, A.R.-T., '20]

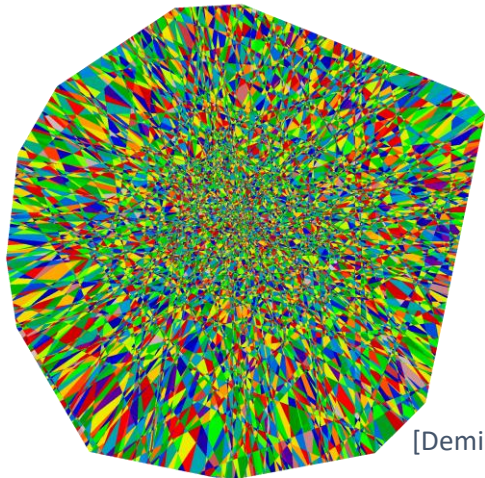
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Et voilà!

We have SUSY AdS with small cosmological constant and all moduli stabilized.

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$$\text{Re}(T_i) = \frac{1}{2} \kappa_{ijk} t^j t^k - \frac{\chi(D_i)}{24} + \frac{1}{(2\pi)^2} \sum_{\mathbf{q} \in \mathcal{M}(X)} q_i \text{GV}_{\mathbf{q}} \text{Li}_2 \left((-1)^{2\mathbf{B} \cdot \mathbf{q}} e^{-2\pi \mathbf{q} \cdot \mathbf{t}} \right)$$

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We must make sure that the vacuum is in the radius of convergence, and that we can find a new point in Kähler moduli space where $D_{T_i} W = 0$ with the corrected volumes.

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With the standard procedure it is impossibly difficult to look deep into rays, so we needed to come up with some tricks.

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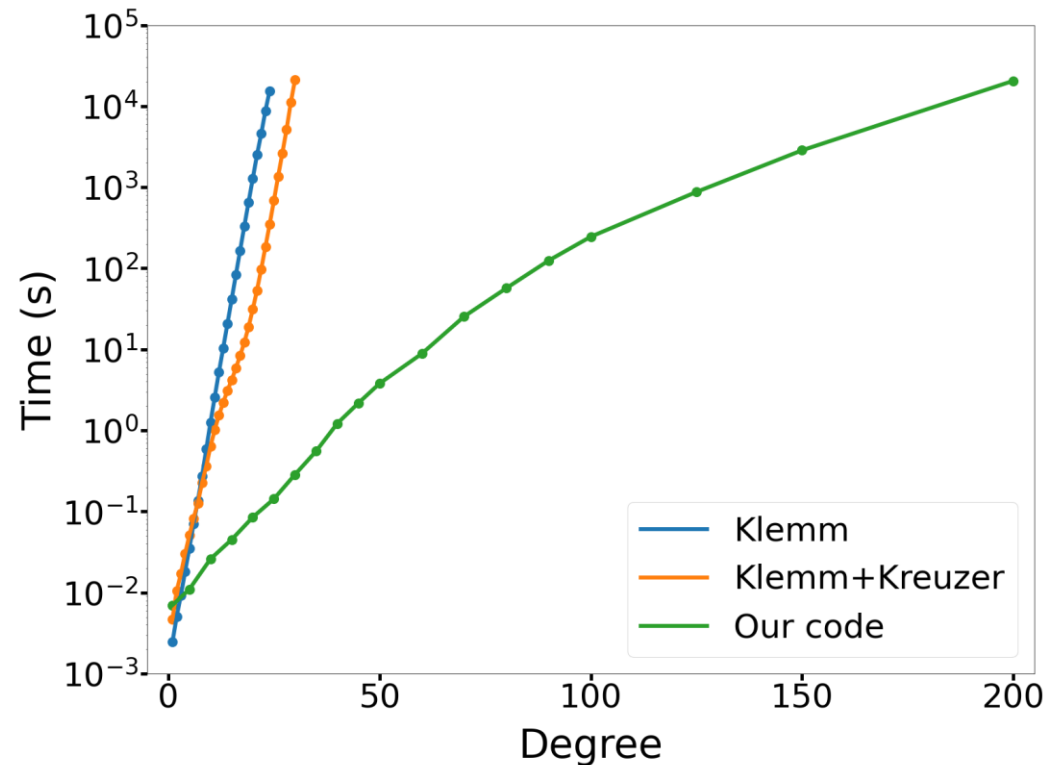
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Here is a comparison for an example at $h^{1,1} = 2$, where the Instanton package can be used.

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We can go all the way up to $h^{1,1} = 491$, and use 100+ million curve classes.

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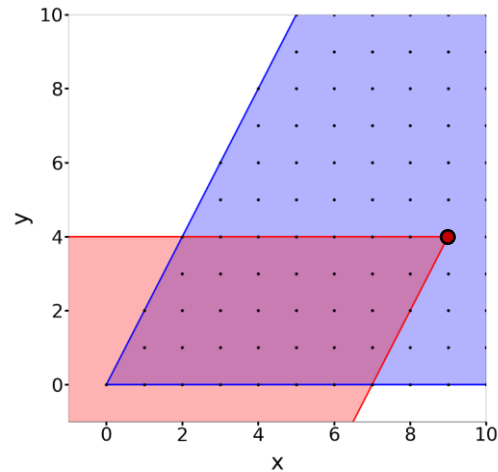
Our examples have $51 \leq h^{1,1} \leq 214$, so we needed to develop new computational tools.

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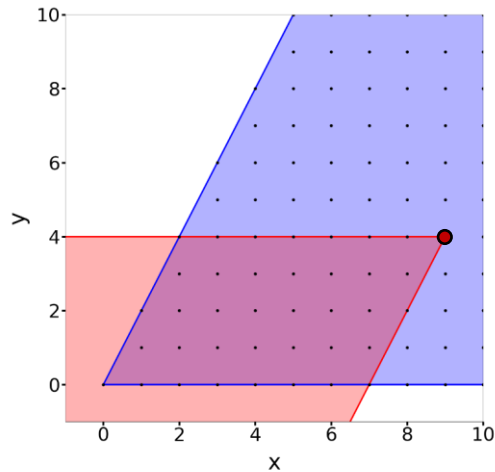
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Blue region is the Mori cone.

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Only curves in the intersection are required for the computation.

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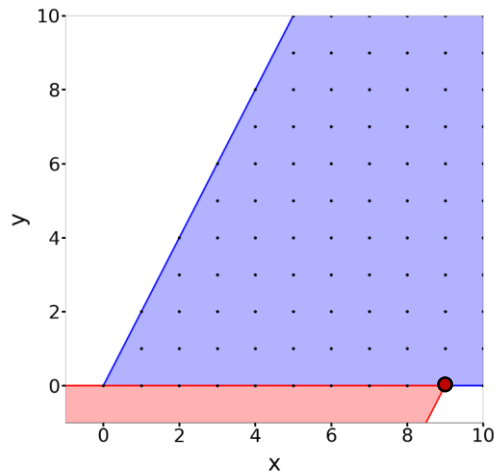


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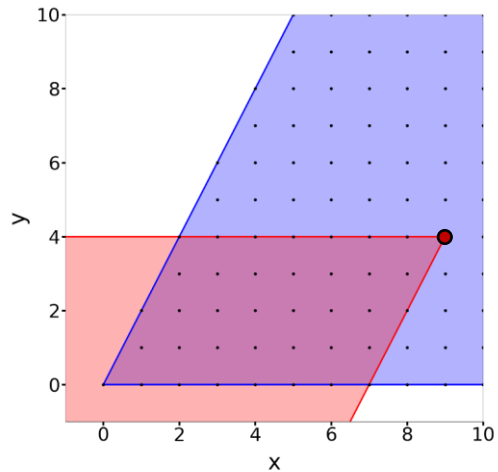
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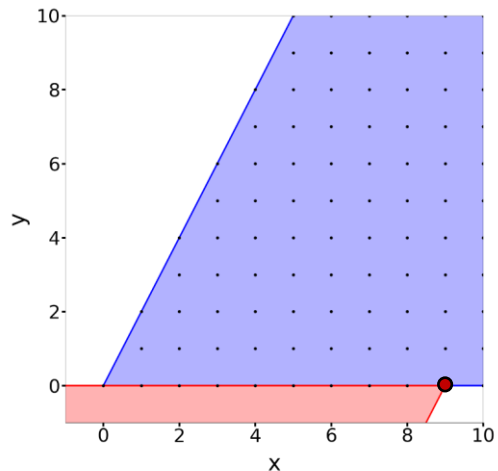


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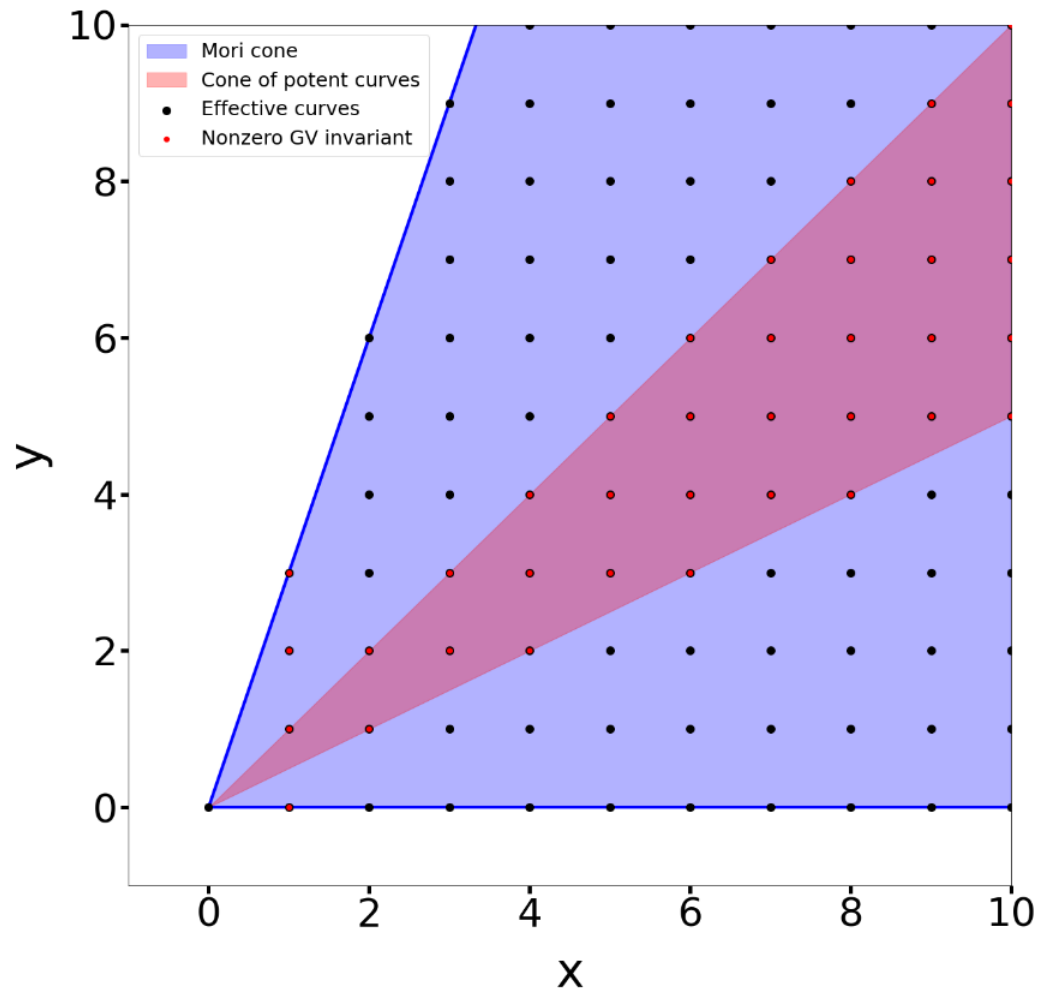


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Computing GV invariants on a d -dimensional face is about as difficult as computing GV invariants for a model with d moduli.

GV invariants are structured

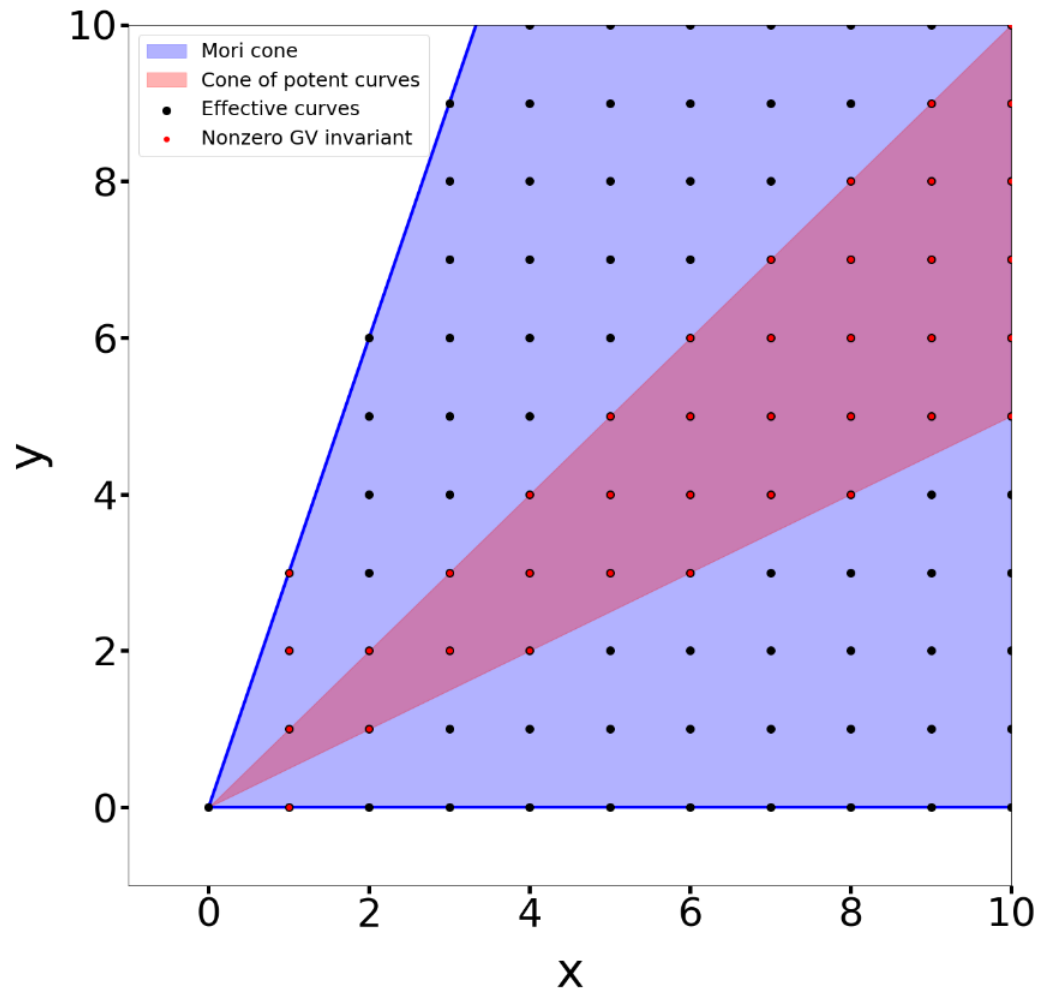
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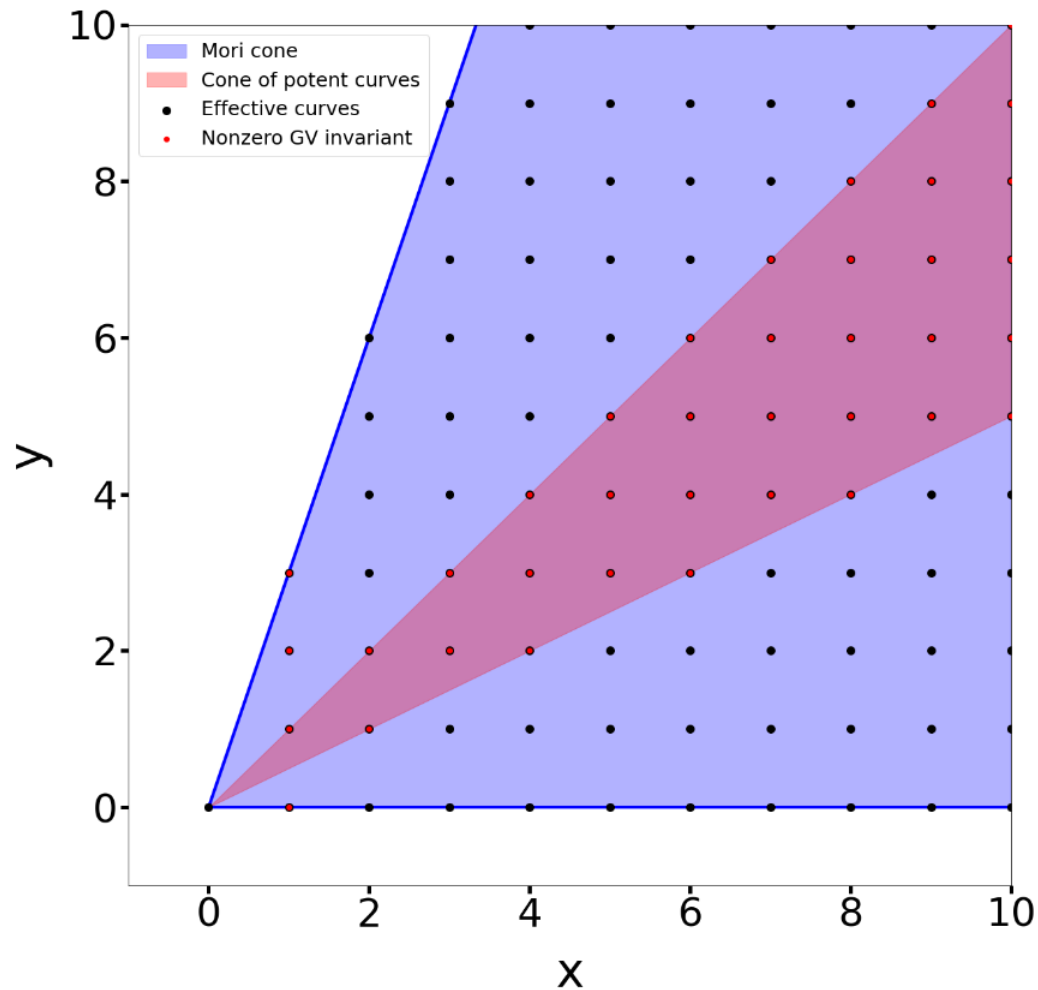


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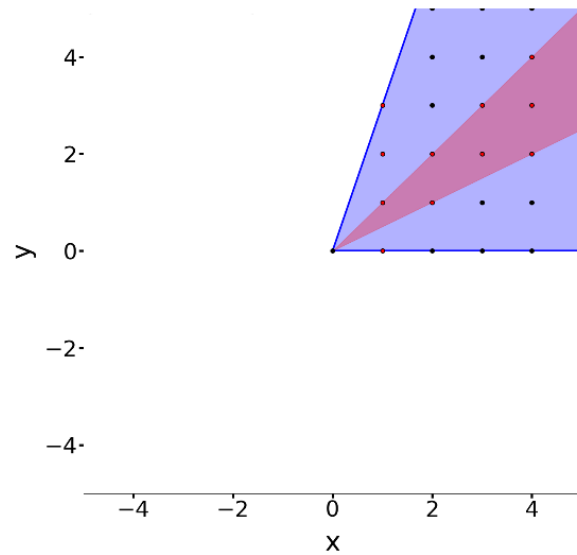
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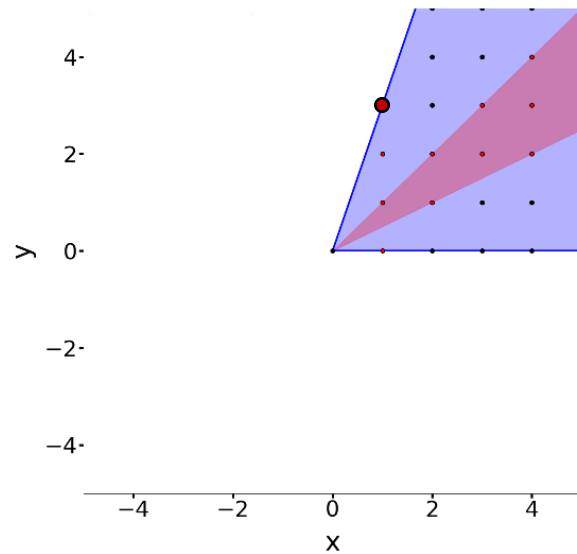
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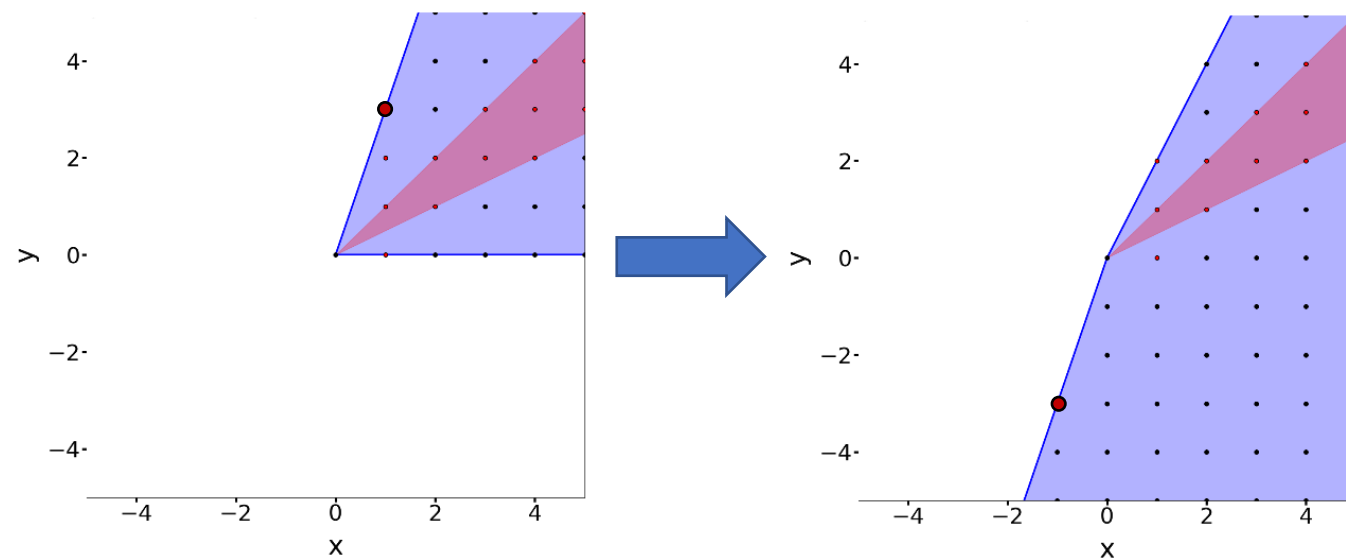
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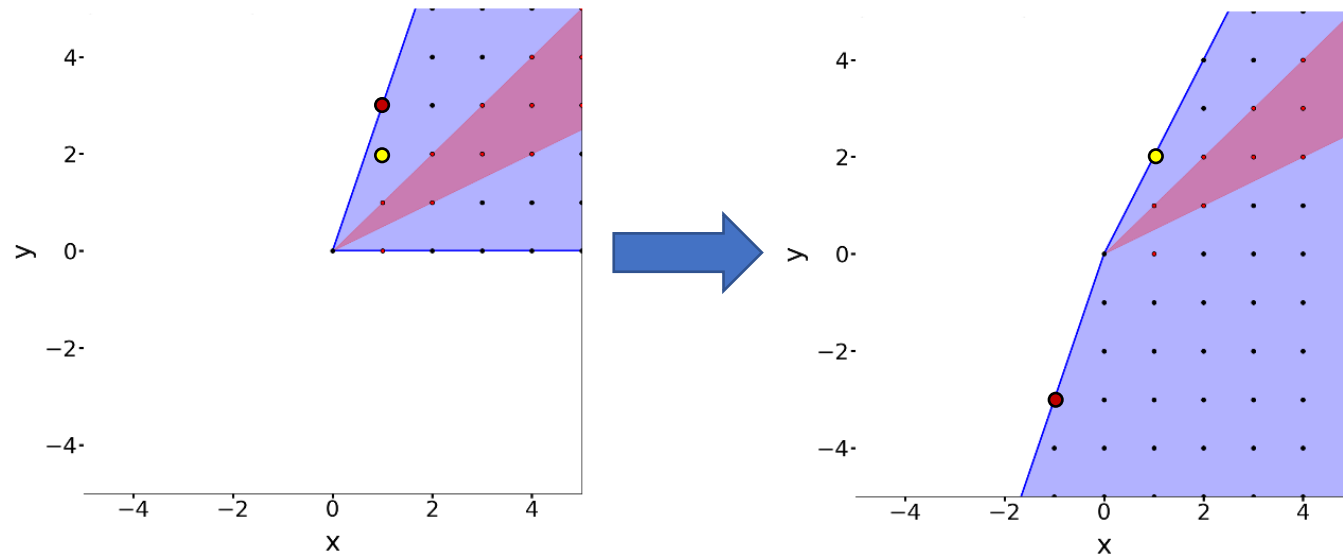
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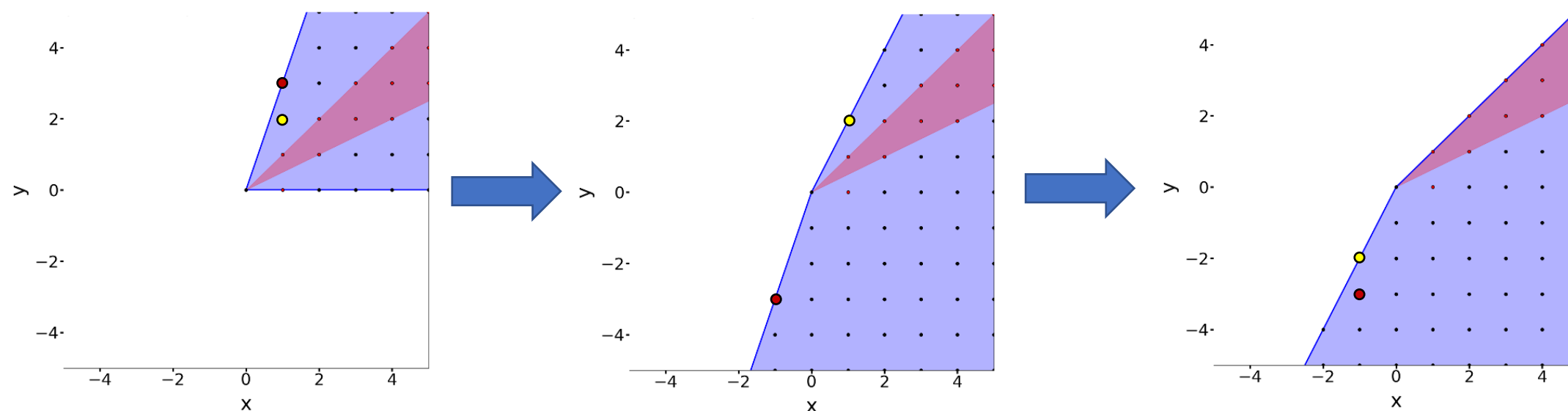
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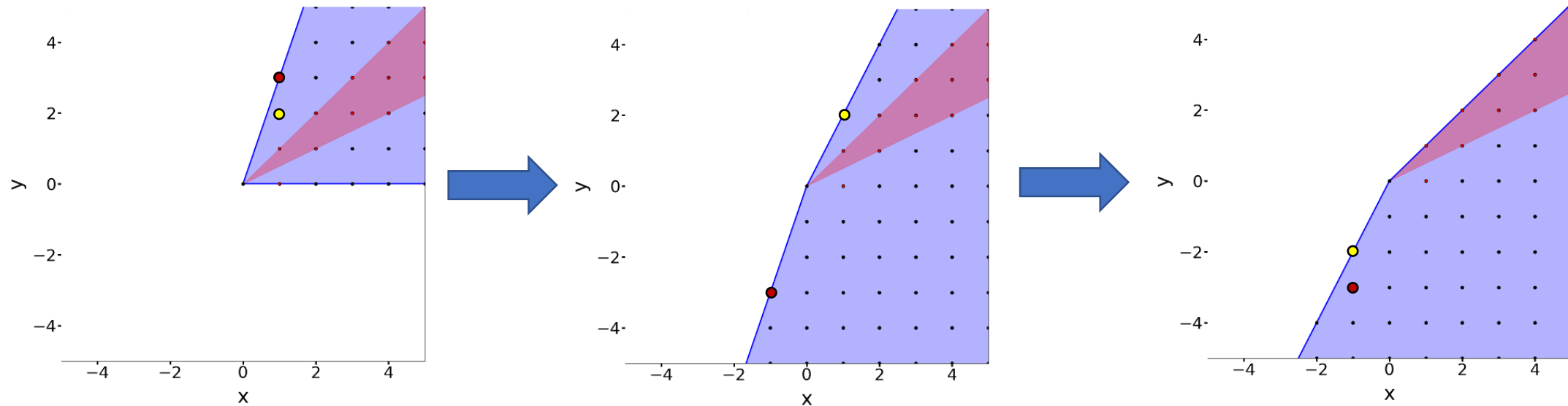
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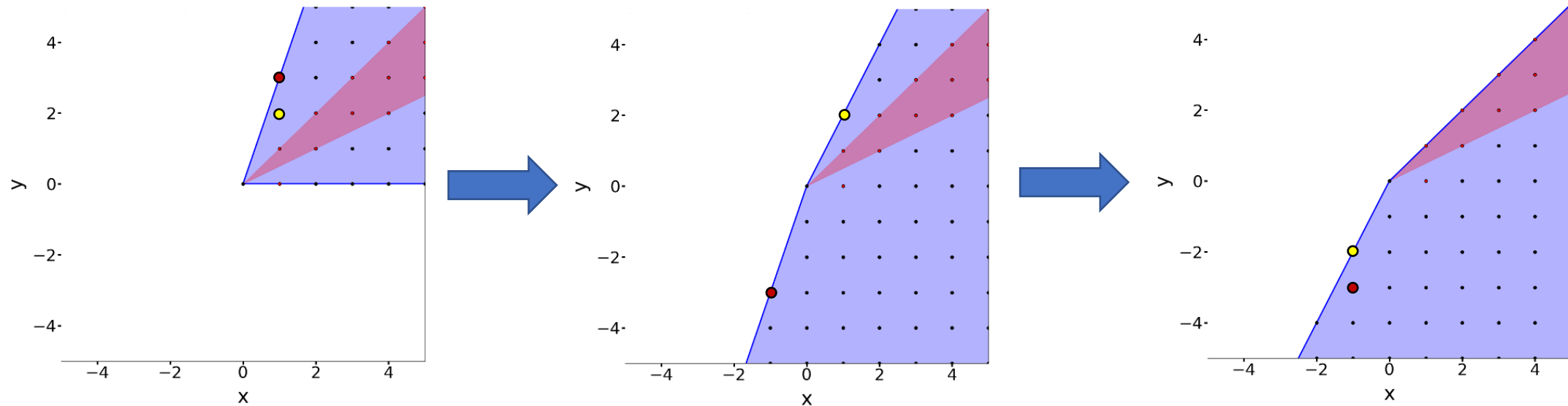
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We can perform flops to go to phases where different parts of the cone of potent rays are in the boundary of the Mori cone.

Then we use the previous trick to compute GV invariants along those faces!

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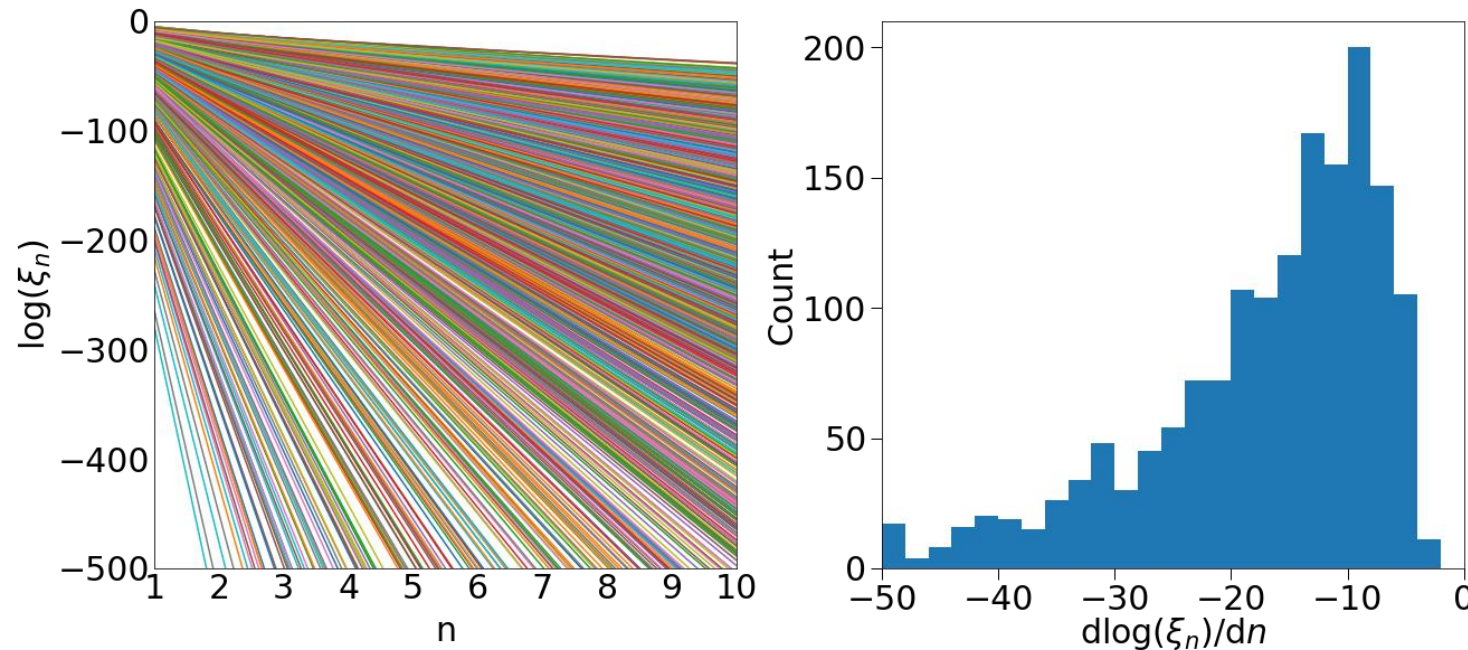
Here we plot the scaling of the contributions along the potent rays. $\xi_n = GV_{n\vec{q}} e^{-2\pi n\vec{q}\cdot\vec{t}}$

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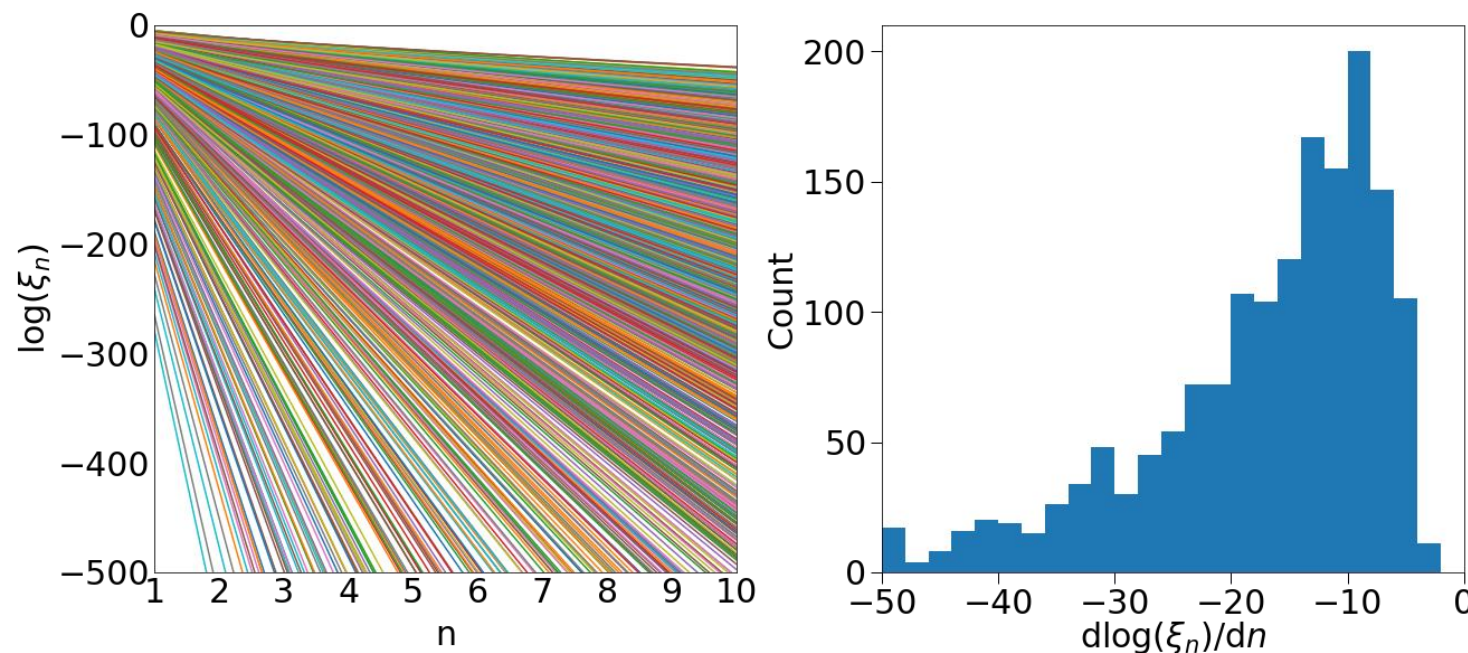


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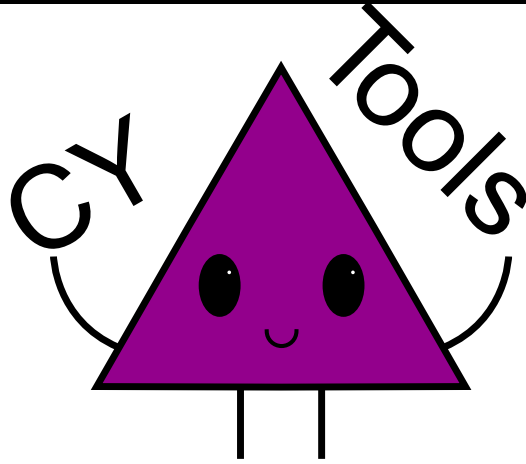
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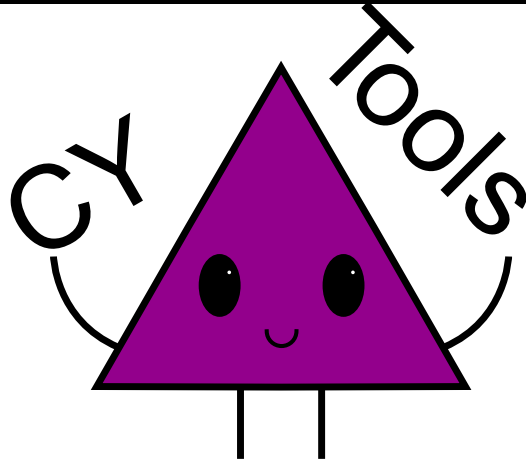
The contributions decay exponentially, so the sum converges!

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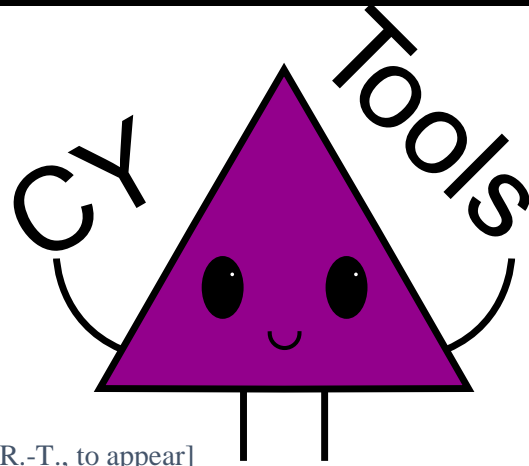
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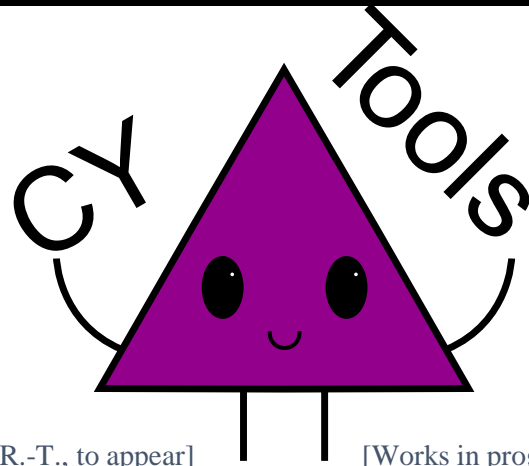
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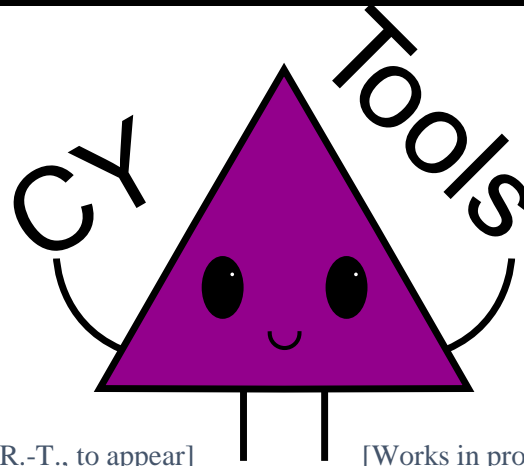


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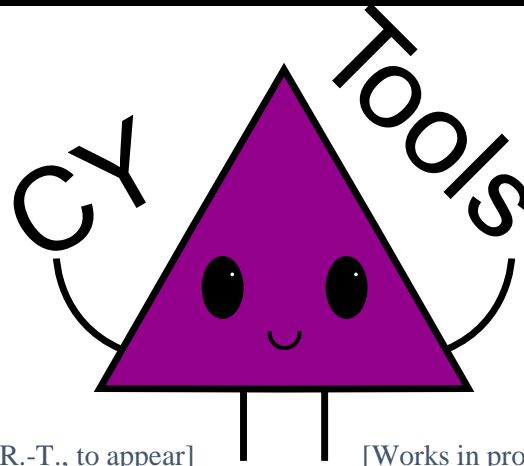


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If any of this sounds interesting to you then let's talk!

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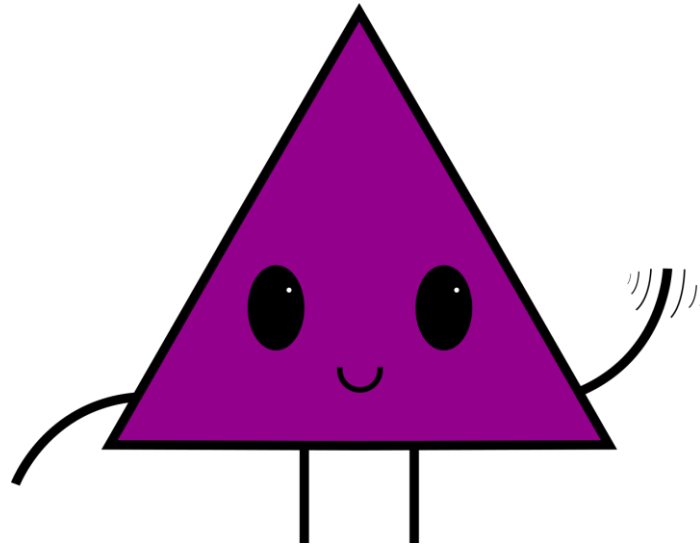
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- Our tools will be included in our CYTools package soon.

Thank you!



Questions?